

## Metal–Metal *d*–*d* Interaction through the Discrete Stacking of Mononuclear M(II) Complexes (M = Pt, Pd, and Cu) within an Organic-Pillared Coordination Cage

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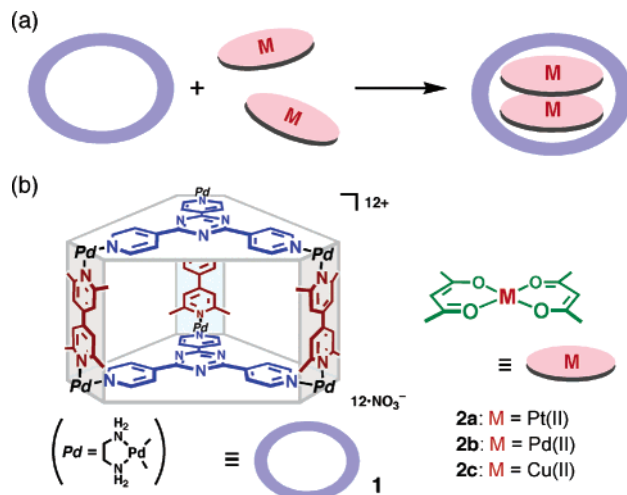
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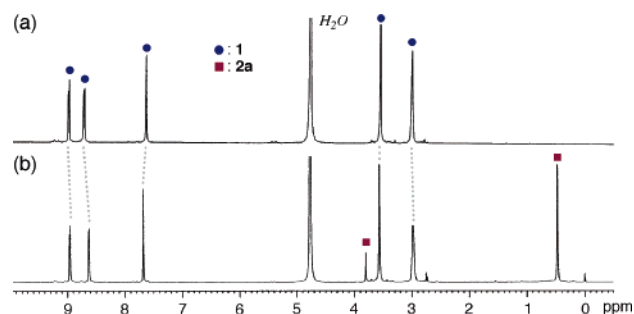
Stacking of planar transition-metal complexes via *d*–*d* interaction between the metal centers promises unique optical, electroconductive, and magnetic properties.<sup>1</sup> Some planar metal complexes, such as Pt(CN)<sub>4</sub><sup>2-</sup> and Pt(NH<sub>3</sub>)<sub>4</sub><sup>2+</sup>/PtCl<sub>4</sub><sup>2-</sup>, assemble into infinitely stacked linear polymers in solid state. Unless bridged by organic ligands, however, common planar metal complexes interact very weakly at the metal centers and seldom form metal–metal bonds via discrete/infinite stacking.<sup>2,3</sup> Here, we report that simple and rather classical metal complexes, M<sup>II</sup>(acac)<sub>2</sub> (M = Pt, Pd, or Cu; acac = acetylacetonato), which have been never shown to form an intermolecular M–M bond,<sup>4</sup> exhibit characteristic metal–metal interaction through accommodation within a coordination cage (Figure 1). The cage we employ here (**1**) has an organic-pillared framework with a large box-shaped hydrophobic cavity, which is ideal to bind two planar molecules as previously reported.<sup>5,6</sup> We show that the M–M interaction is clearly demonstrated by spectroscopies, crystallographic analysis (for M = Pt(II)), and electron spin–spin coupling (for M = Cu(II)).

From cage **1** and Pt(acac)<sub>2</sub> (**2a**), inclusion complex **1**⊃(**2a**)<sub>2</sub> was quantitatively obtained by a very simple procedure. Complex **2a** (4 molar equiv) was suspended in a D<sub>2</sub>O solution of **1** (10 mM), and the mixture was stirred at room temperature for 30 min. After the removal of excess **2** by filtration, <sup>1</sup>H NMR analysis of the clear solution revealed the quantitative formation of **1**⊃(**2a**)<sub>2</sub> complex. Highly upfield-shifted signals of **2a** were observed at 3.80 and 0.47 ppm, indicating the enclathration of **2a** within **1** (Figure 2b). The host–guest ratio (1:2) was estimated by the integral ratio of **1** and **2a**. CSI–MS measurement also showed the formation of **1**⊃(**2a**)<sub>2</sub> very clearly. Intense signals were observed at *m/z* 642.9, 740.2, and 904.4, which correspond to [**1**⊃(**2a**)<sub>2</sub> – 6·NO<sub>3</sub><sup>-</sup> + 6·DMF]<sup>6+</sup>, [**1**⊃(**2a**)<sub>2</sub> – 5·NO<sub>3</sub><sup>-</sup> + 3·DMF]<sup>5+</sup>, and [**1**⊃(**2a**)<sub>2</sub> – 4·NO<sub>3</sub><sup>-</sup> + DMF]<sup>4+</sup>, respectively. The **1**⊃(**2a**)<sub>2</sub> host–guest complex was shown to be stable because guest dissociation from cage **1** under the CSI–MS conditions was hardly observed.

Pt(II)–Pt(II) interaction in **1**⊃(**2a**)<sub>2</sub> was displayed by X-ray crystallographic analysis (Figure 3). Orange-colored single crystals suitable for X-ray analysis grew by the slow evaporation of water from an aqueous solution of **1**⊃(**2a**)<sub>2</sub> at room temperature for 3 days.<sup>7</sup> The crystal structure reveals the stacking of two molecules of **2a** in such a way that two Pt(II) metals exist on the C<sub>2</sub> symmetry axis of the stacked dimer. The Pt(II)–Pt(II) distance is 3.32 Å, being characteristic to typical Pt(II)–Pt(II) *d*–*d* interaction (ca. <3.5 Å).<sup>1–3</sup> Around the Pt(II)–Pt(II) axis, the two Pt(acac)<sub>2</sub> molecules are twisted by 24.8° to reduce the steric repulsion



**Figure 1.** (a) Discrete stacking of planar metal complexes within a cage. (b) Chemical structures of organic-pillared coordination cage **1** and M<sup>II</sup>(acac)<sub>2</sub> complexes **2**.



**Figure 2.** NMR spectra (500 MHz, D<sub>2</sub>O, rt) of (a) **1** and (b) **1**⊃(**2a**)<sub>2</sub>.

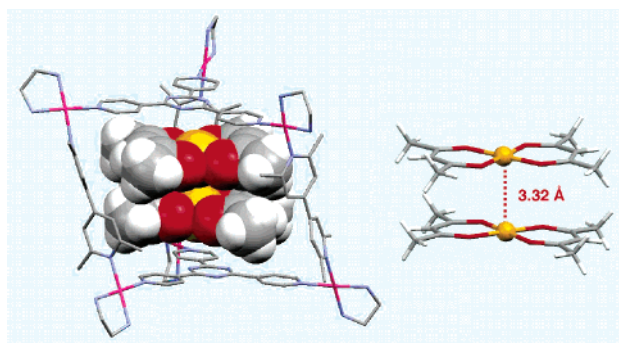
between the methyl groups. Cage **1** itself is also twisted by 23.5° to maximize the host–guest interaction.

UV–vis measurement of powdered **1**⊃(**2a**)<sub>2</sub> also evidenced Pt(II)–Pt(II) interaction. The spectrum showed absorption band around 500 nm (Figure 4), which is featured by *d*–*d* interaction of Pt(II) complexes.<sup>1–3</sup> For **1** and **2a**, no absorption was observed above 450 nm. The orange color of the complex turned almost colorless when the complex was dissolved in water, suggesting that the Pt(II)–Pt(II) interaction became weaker in solution. However, <sup>195</sup>Pt NMR of **1**⊃(**2a**)<sub>2</sub> showed a broad signal at –280 ppm, which was considerably downfield-shifted ( $\Delta\delta = 130$  ppm) compared with that of free **2a**.<sup>8,9</sup> This indicated that the Pt(II)–Pt(II) interaction still remained to some extent even in solution.<sup>2,3</sup>

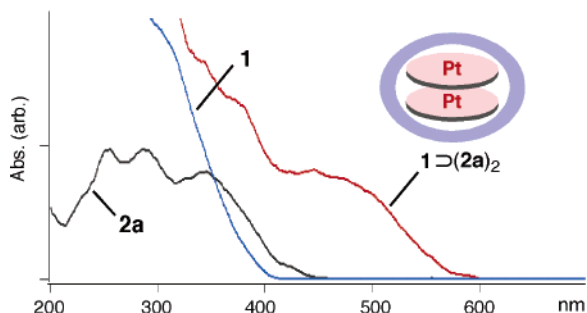
We also examined the accommodation of analogous Pd(acac)<sub>2</sub> (**2b**) by cage **1**. Again, **1**⊃(**2b**)<sub>2</sub> complex was obtained quantita-

<sup>†</sup> The University of Tokyo and CREST, Japan Science and Technology Agency (JST).

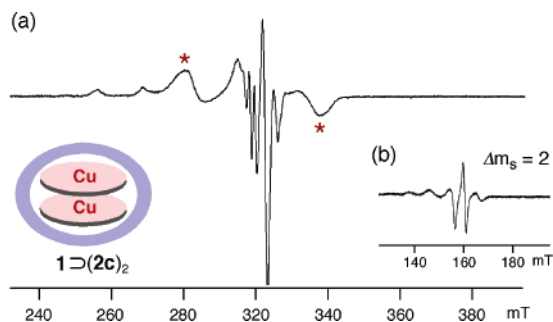
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**Figure 3.** X-ray structure of  $1\supset(2a)_2$  (left) and selected structure of  $2a$  (yellow = Pt, red = oxygen) within  $1$  (right).



**Figure 4.** UV-vis spectra of  $1$ ,  $2a$ , and  $1\supset(2a)_2$  in solid state.



**Figure 5.** ESR spectra ( $H_2O$ , 103 K,  $MnO$  as external standard) of (a)  $1\supset(2c)_2$  and (b) the forbidden transition.

tively, and it exhibited a new absorption band around 450 nm in solid state, which is attributed to Pd(II)–Pd(II) interaction.<sup>2,9</sup>

For a  $d^9$  square-planar metal complex,  $Cu(acac)_2$  ( $2c$ ), we expected spin–spin interaction between two Cu(II) centers by spontaneous stacking of  $2c$  within  $1$ . In fact, the treatment of  $2c$  with an aqueous solution of  $1$  gave  $1\supset(2c)_2$  complex, and broad signals in ESR at low temperature (labeled by the asterisks (\*) in Figure 5a) grew up with the formation of  $1\supset(2c)_2$  complex.<sup>9</sup> The broad signals are attributed to the  $d$ – $d$  interaction between the copper centers of  $1\supset(2c)_2$  in  $S = 1$  state, and the large splitting of ca. 60 mT is due to the spin–spin interaction. The presence of a triplet species ( $S = 1$  state) is confirmed by observing  $\Delta m_s = 2$  (forbidden transition) signal at 160 mT (half of 320 mT for  $\Delta m_s = 1$  transition) (Figure 5b). The splitting of 60 mT gives an averaged interspin distance of 3.6 Å.<sup>10–12</sup> A sharply split signal around 325 mT is assigned as a doublet radical ( $S = 1/2$ ), which suggests the presence of  $1\supset 2c$  as a minor component.<sup>9</sup> It should be noted that

the estimated Cu(II)–Cu(II) distance is longer than the Pt(II)–Pt(II) distance determined by X-ray analysis. This comes from the different fashion of  $d$ – $d$  interaction between the two cases. Whereas the Pt(II)–Pt(II) interaction involves the bonding overlap of linearly disposed  $d_{z^2}$  and  $p_z$  orbitals, the Cu(II)–Cu(II) interaction involves nonbonding stacking of two  $d_{x^2-y^2}$  orbitals.

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**Supporting Information Available:** Experimental details and spectroscopic data (PDF). This material is available free of charge via the Internet at <http://pubs.acs.org>.

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- (7) X-ray crystal data of  $1\supset(2)$ :  $C_{10}H_{148}N_{42}O_{80}Pd_6Pt_2$ ,  $M = 4511.28$ , crystal dimensions  $0.30 \times 0.20 \times 0.20$  mm<sup>3</sup>, monoclinic space group  $P2_1/c$ ,  $a = 18.682(3)$  Å,  $b = 36.535(6)$  Å,  $c = 26.576(4)$  Å,  $\beta = 108.157(3)^\circ$ ,  $V = 17236(5)$  Å<sup>3</sup>,  $Z = 4$ ,  $\rho_{\text{calcd}} = 1.738$  g cm<sup>-3</sup>,  $F(000) = 8984$ , radiation,  $\lambda(\text{Mo K}\alpha) = 0.71073$  Å,  $T = 80(2)$  K, reflections collected/unique 194260/41116 ( $R_{\text{int}} = 0.3143$ ). The structure was solved by direct methods (SHELXL-97) and refined by full-matrix least-squares methods on  $F^2$  with 1753 parameters.  $R_1 = 0.1496$  ( $I > 2\sigma(I)$ ),  $wR_2 = 0.2894$ , GOF 1.025; max/min residual density 3.518/–8.602 e Å<sup>-3</sup>. CCDC reference number 270934.
- (8) A control experiment was done with pyrazine-pillared cage  $1'$  (Kumazawa, K.; Biradha, K.; Kusakawa, T.; Okano, T.; Fujita, M. *Angew. Chem., Int. Ed.* **2003**, *42*, 3909–3913). In contrast to  $1\supset(2a)_2$ ,  $1'\supset 2a$  complex showed an upfield-shifted signal in <sup>195</sup>Pt NMR compared with free  $2a$ . Thus, significant downfield shift in  $1\supset(2a)_2$  is due to the influence of the cage but to the Pt(II)–Pt(II) interaction.
- (9) See Supporting Information.
- (10) Point dipole approximation was applied:

$$D = \pi \frac{3}{2} \left( \frac{\mu_0}{4\pi} \right) \frac{(g\beta)^2}{R^3} = (2.80 \times 10^3) \frac{1}{R^3}$$

where  $D$  (mT) is the splitting of the dipole interaction and  $R$  (Å) is the averaged interspin distance.

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